

Fuerzas Centrales

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3. FUERZAS CENTRALES

$$\vec{F}(\vec{r}) = F(r)\hat{r}$$

Problema de 2 cuerpos: Si $\vec{F}_{12} = -\vec{F}_{21}$ i $\frac{\vec{F}_1}{m_1} = \frac{\vec{F}_2}{m_2}$

$$\begin{cases} m_1\ddot{\vec{r}}_1 = \vec{F}_{21} + \vec{F}_1 \\ m_2\ddot{\vec{r}}_2 = \vec{F}_{12} + \vec{F}_2 \end{cases} \mapsto \begin{cases} M\ddot{\vec{R}} = \vec{F} \\ \mu\ddot{\vec{r}} = \vec{F}_{21} \end{cases}, \text{ donde:}$$

$$M = m_1 + m_2, \quad \vec{F} = \vec{F}_1 + \vec{F}_2, \quad \vec{r} = \vec{r}_1 - \vec{r}_2, \\ \vec{R} = \frac{m_1\vec{r}_1 + m_2\vec{r}_2}{m_1 + m_2}, \quad \mu = \frac{m_1m_2}{m_1 + m_2}$$

Dinámica en \mathbb{R}^3

$$1. \vec{p} = m\vec{v}; \quad \vec{I}_{1,2} = \int_{t_1}^{t_2} \vec{F}(\vec{r}, \dot{\vec{r}}, t) dt \rightarrow \vec{F} = \frac{d\vec{p}}{dt}$$

$$2. \left. \begin{aligned} \vec{L} &= \vec{r} \times \vec{p} \\ \vec{N} &= \vec{r} \times \vec{F} \end{aligned} \right\} \mapsto \frac{d\vec{L}}{dt} = \vec{N} \rightarrow \Delta\vec{L} = \int_{t_1}^{t_2} \vec{N} dt$$

$$3. T = \frac{1}{2}m\vec{v}^2; \quad W = \int_{\vec{r}_1}^{\vec{r}_2} \vec{F}(\vec{r}, \dot{\vec{r}}, t) \cdot d\vec{r}; \quad \frac{dT}{dt} = \vec{F} \cdot \vec{v}$$

4. Sistemas conservativos: $\vec{F} = \vec{F}(r) \quad \vec{\nabla} \times \vec{F} = \vec{0}$

$$U(\vec{r}) = - \int_{\vec{r}_R}^{\vec{r}} \vec{F}(\vec{r}') \cdot d\vec{r}' \rightarrow \vec{F}(\vec{r}) = -\vec{\nabla}U(\vec{r}) \\ E = \frac{1}{2}m\vec{v}^2 + U(\vec{r}) = \mathcal{C}$$

Propiedades de los campos centrales

$$\vec{F}(\vec{r}) = F(r)\hat{r} = \frac{x}{r}F(r)\hat{i} + \frac{y}{r}F(r)\hat{j} + \frac{z}{r}F(r)\hat{k}$$

$$\vec{L} = \vec{C}, \quad \Delta E = 0, \quad \vec{\nabla} \times F(r)\hat{r} = \vec{0}$$

$$U(r) = - \int_{r_s}^r F(r')dr' \quad \vec{F} = -\vec{\nabla}U(\vec{r}); \quad F(r) = -\frac{dU}{dr}$$

$$E = \frac{1}{2}m\vec{v}^2 + U(r)$$

Métodos de resolución

$$x = r \cos \varphi, \quad y = r \sin \varphi, \quad r = \sqrt{x^2 + y^2}, \quad \varphi = \arctan \frac{y}{x}$$

$$\hat{r} = \cos \varphi \hat{i} + \sin \varphi \hat{j}, \quad \hat{\varphi} = -\sin \varphi \hat{i} + \cos \varphi \hat{j} \\ \dot{\hat{r}} = \dot{\varphi} \hat{\varphi}, \quad \dot{\hat{\varphi}} = -\dot{\varphi} \hat{r}$$

1. Aplicación de la segunda ley de Newton

$$\vec{r} = r\hat{r}$$

$$\dot{\vec{r}} = \dot{r}\hat{r} + r\dot{\varphi}\hat{\varphi}$$

$$\ddot{\vec{r}} = (\ddot{r} - r\dot{\varphi}^2)\hat{r} + (2\dot{r}\dot{\varphi} + r\ddot{\varphi})\hat{\varphi}$$

$$m\ddot{\vec{r}} = \vec{F}(\vec{r}) = F(r)\hat{r} + 0 \cdot \hat{\varphi} \rightarrow \begin{cases} m(\ddot{r} - r\dot{\varphi}^2) = F(r) \\ 2\dot{r}\dot{\varphi} + r\ddot{\varphi} = 0 \end{cases}$$

$$L = mr^2\dot{\varphi} = \mathcal{C}$$

$$\begin{cases} m\ddot{r} = F(r) + \frac{L}{mr^3} \equiv F'(r) \\ \dot{\varphi} = \frac{L}{mr^2} \end{cases}$$

2. Teoremas de conservación

$$\begin{cases} L = mr^2\dot{\varphi} \\ E = \frac{1}{2}m\vec{v}^2 + U(\vec{r}) \end{cases} \rightarrow \begin{cases} L = mr^2\dot{\varphi} \\ E = \frac{1}{2}m\dot{r}^2 + \frac{1}{2}mr^2\dot{\varphi}^2 + U(r) \end{cases}$$

$$\dot{\varphi}^2 = \frac{L^2}{m^2r^4}$$

$$E = \frac{1}{2}m\dot{r}^2 + \frac{1}{2}\frac{L^2}{mr^2} + U(r) = \frac{1}{2}m\dot{r}^2 + U'(r)$$

$$\blacksquare \int_{r_0}^r \frac{dr'}{\sqrt{\frac{2}{m}[E - U'(r')]} = \int_0^t dt' \mapsto r(t) = \Phi^{-1} \left(\sqrt{\frac{2}{m}}t \right)$$

$$\blacksquare \varphi(t) = \int_0^t \frac{L}{mr^2(t')} dt'$$

3. Análisis semicuantitativo de $U'(r)$: $E \geq U'(r)$

$$\blacksquare \text{Puntos de retorno: } \left(\frac{dU'}{dr} \right)_{r_R} \neq 0; \quad E = U'(r_R)$$

\blacksquare Puntos de equilibrio:

$$\left(\frac{dU'}{dr} \right)_{r_E} = 0 \rightarrow \begin{cases} \left(\frac{d^2U'}{dr^2} \right)_{r_E} > 0 \rightarrow \text{estable} \\ \left(\frac{d^2U'}{dr^2} \right)_{r_E} < 0 \rightarrow \text{inestable} \end{cases}$$

4. Ecuación de la trayectoria

$$r = r(t)\varphi = \varphi(t) \rightarrow r = r(\varphi) \quad u(\varphi) \equiv \frac{1}{r(\varphi)}$$

$$\frac{d^2u}{d\varphi^2} + \frac{m}{L^2u^2}F\left(\frac{1}{u}\right) + u = 0$$

Fuerza del tipo $F(r) = \frac{k}{r^2}$

Estudio semicuantitativo usando $U'(r)$

$$U'(r) = U(r) + \frac{L^2}{2mr^2} \mapsto U'(r) = \frac{k}{r} + \frac{L^2}{2mr^2}$$

1. $k > 0$ fuerza repulsiva; No extremos, no r_E , sólo 1 r_R .

$$\frac{1}{r_1} = -\frac{mK}{L^2} + \sqrt{\left(\frac{mk}{L^2}\right)^2 + \frac{2mE}{L^2}} > 0 \rightarrow \text{Órbita ilimitada abierta}$$

2. $k < 0$ $E \geq U_0$

$$r_0 = -\frac{L^2}{mk} \quad U_0 = U'(r_0) = -\frac{mk^2}{2L^2}$$

- $E = U_0$

$$\text{M.R.U} \rightarrow \begin{cases} r_0 = -\frac{L^2}{mk} \\ \dot{\varphi}_0 = \frac{L}{mr_0^2} \end{cases}$$

- $U_0 < E < 0$

$$\frac{1}{r_{1,2}} = -\frac{mk}{L^2} \pm \sqrt{\left(\frac{mk}{L^2}\right)^2 - \frac{2m|E|}{L^2}}$$

- $E \geq 0$

$$\frac{1}{r_1} = -\frac{mk}{L^2} + \sqrt{\left(\frac{mk}{L^2}\right)^2 + \frac{2mE}{L^2}}$$

Ecuación de la trayectoria

$$u(\varphi) = A \cos \varphi - \frac{mk}{L^2}$$

Determinar $A \rightarrow \varphi = 0$ ó $\varphi = \pi$:

$$A = \sqrt{\left(\frac{mk}{L^2}\right)^2 + \frac{2mE}{L^2}}$$

$$\boxed{r(\varphi) = \frac{1}{A \cos \varphi + B}} \quad B = -\frac{mk}{L^2}$$

1. ELIPSE: $E < 0$, $k < 0$ $\varepsilon < 1$

$$r(\varphi) = \frac{a(1 - \varepsilon^2)}{1 + \varepsilon \cos \varphi} = \frac{1}{B + A \cos \varphi}$$

2. HIPÉRBOLA: $E > 0$ si $k > 0$ rama -, si $k < 0$ rama +, $\varepsilon > 1$

$$r(\varphi) = \frac{a(\varepsilon^2 - 1)}{\pm 1 + \varepsilon \cos \varphi} = \frac{1}{B + A \cos \varphi}$$

$$\varepsilon = \frac{A}{|B|} \quad a = \left| \frac{B}{A^2 - B^2} \right|$$

3. PARÁBOLA: $E = 0$; $k < 0$; $\varepsilon = 1$

$$r(\varphi) = \frac{a}{1 + \cos \varphi}$$

$$\varepsilon = \sqrt{1 + \frac{2EL^2}{mk^2}} \quad a = \left| \frac{k}{2E} \right|$$

Problema de Kepler

1. Los planetas siguen órbitas elípticas con el sol en uno de sus focos.

2. La velocidad aerolar es constante. $\frac{dS}{dt} = \frac{L}{2m}$

3. $T^2 \propto a^3$ $T = \frac{4\pi^2}{GM} a^3$